

LETTER TO THE EDITOR

Finding short GRB remnants in globular clusters: the VHE gamma-ray source in Terzan 5

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ABSTRACT

Context. Globular cluster are believed to boost the rate of compact binary mergers that may launch a certain type of cosmological gamma-ray burst (GRB). Therefore globular clusters appear to be potential sites to search for remnants of such GRBs.

Aims. The very-high-energy (VHE) gamma-ray source HESS J1747-248 recently discovered in the direction of the Galactic globular cluster Terzan 5 is investigated for being a GRB remnant.

Methods. Signatures created by the ultrarelativistic outflow, the subrelativistic ejecta and the ionizing radiation of a short GRB are estimated for an expected age of such a remnant of $t \gtrsim 10^4$ years.

Results. The kinetic energy of a short GRB could roughly be adequate for powering the VHE source in a hadronic scenario. The age of the proposed remnant estimated from its extension possibly agrees with the occurrence of such events in the Galaxy. Subrelativistic merger ejecta could shock-heat the ambient medium.

Conclusions. Further VHE observations can probe for the presence of a break towards lower energies expected for particle acceleration in ultrarelativistic shocks. Deep X-ray observations would have the potential to examine whether there is any thermal plasma heated by the subrelativistic ejecta. The identification of a GRB remnant in our own Galaxy may also help to explore the effect of such a highly energetic event on the Earth.

Key words. ISM: supernova remnants – Gamma-ray burst: general – Galaxy:globular clusters: individual: Terzan 5

1. Introduction

Gamma-ray bursts (GRBs) are enigmatic explosions, detected usually at cosmological distances. They appear to come in two distinct flavors according to their duration (see e.g. Gehrels et al., 2009, for a review). Powerful explosions usually leave long-living remnants behind that potentially could be studied in the local Universe as relics of the original events.

Long bursts (duration longer than about 2 s) are generally considered to be launched by the death of a massive star and several structures in our Galaxy have been suggested to be the remnants of such events. These proposed remnants are in many cases linked to sources of very-high-energy (VHE, >100 GeV) gamma-ray emission. The VHE gamma-ray source HESS J1303-631 (Aharonian et al., 2005) has been suggested as the remnant of a long GRB that happened $\gtrsim 10^4$ years ago (Atoyan et al., 2006). Also the structure W49B has been suspected to be a remnant of a GRB (Ioka et al., 2004). Recently it has been argued that the population of unidentified TeV sources (see e.g. Aharonian et al., 2008) may be dominated by GRB/Hypernova remnants (Ioka & Mészáros, 2010).

For short GRBs (duration shorter than 2 s), compact binary mergers as central engine have been identified as the preferred scenario (e.g. Gehrels et al., 2005; Lee et al., 2005). Such compact binaries may be efficiently created in globular clusters since these environments feature high densities of very old stars in their cores. Indeed in the globular cluster M 15 already a close system that consists of two neutron stars (NS) has been discovered (Anderson et al., 1990). Based on these facts, it has been

argued that a considerable fraction of all NSNS binaries are formed in globular clusters (Grindlay et al., 2006). It has been claimed that the rate of short bursts in the local Universe is dominated by mergers of dynamically formed compact binaries in globular clusters (Salvatera et al., 2008; Guetta & Stella, 2009). Therefore, globular clusters might be the prime environment for seeking remnants of short GRBs. Potential signatures of remnants of compact binary mergers have already been discussed (see Domainko & Ruffert, 2005, 2008).

The H.E.S.S. collaboration has very recently reported the detection of a VHE gamma-ray source HESS J1747-248 in the direction of the galactic globular cluster Terzan 5 (Abramowski et al., 2011) with a flux above 440 GeV of $(1.2 \pm 0.3) \times 10^{-12} \text{ cm}^{-2} \text{ s}^{-1}$. The source is extended with intrinsic extension of $9'.6 \pm 2'.4$, and there appears to be an offset from the cluster core position by about $4'.0 \pm 1'.9$. Terzan 5 is located at a distance of 5.9 kpc (Ferraro et al., 2009), at RA(J2000) $17^{\text{h}}48^{\text{m}}04^{\text{s}}.85$ and Dec $-24^{\circ}46'44''.6$ ($l 3.8^{\circ}, b 1.7^{\circ}$), and exhibits a core radius of $0.15'$ a half-mass radius of $0.52'$ and a tidal radius of $4.6'$ (Lanzoni et al., 2010). Terzan 5 is the globular cluster with the highest expected rates of close stellar encounters (Pooley & Hut, 2006) and with the largest population of millisecond pulsars (33) discovered up to now (Ransom, 2008). Gamma rays of likely magnetospheric origin produced by millisecond pulsars have been detected in the GeV range in Terzan 5 (Kong et al., 2010; Abdo et al., 2010a). Diffuse X-ray emission extending beyond the half-mass radius has also been reported (Eger et al., 2010). In the vicinity of Terzan 5 several structures in the radio band have been found (Clapson et al., 2011). Terzan 5 is a predicted VHE gamma-ray emitter where energetic electrons produced by the large population of millisecond

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ond pulsars up-scatter stellar photons to gamma-ray energies (Venter et al., 2009; Bednarek & Sitarek, 2007). In contrast to these models here VHE gamma-ray emission originating from collisions of hadronic cosmic rays with ambient target nuclei and subsequent π^0 decay is explored. A short GRB is adopted as the accelerator of the cosmic rays.

The paper is organized as follows: In Sect. 2 potential indications for a hadronic VHE gamma-ray production in Terzan 5 are discussed, in Sect. 3 signatures left behind by the ultrarelativistic outflow of an ancient GRB in Terzan 5 are assessed, in Sect. 4 signatures related to the subrelativistic ejecta are explored, and in Sect. 5 potential traces of ionizing radiation emitted by the GRB are investigated.

2. Potential indications for a hadronic scenario

Some of the properties of the VHE gamma-ray source that were detected in the direction of Terzan 5 appear to challenge leptonic models for the gamma-ray production. In particular, the extension, indication for an offset of the source from the globular cluster core and a power-law spectrum are not self-evident for such a scenario. The intensity of leptonic inverse-Compton (IC) radiation scales linearly with the energy density of the target photon field. For Terzan 5 the energy density of the stellar photon field drops from about 1000 eV/cm^{-3} in the core region to 40 eV/cm^{-3} at the half-mass radius (Venter et al., 2009) to about a few eV/cm^{-3} at the extension of the VHE source. Consequently a very centrally peaked source centered on the GC would be expected, which does not seem to be supported by the H.E.S.S. observations. IC emission in the VHE range should be accompanied by synchrotron emission in the X-ray band. Diffuse X-ray emission centered on Terzan 5 of possible synchrotron origin has indeed been discovered (Eger et al., 2010) but the potential offset of the VHE gamma-ray emission from the peaks of the X-ray emission and radiation field challenges a leptonic scenario. Furthermore, since the optical to near-infrared starlight photon field should be up-scattered by the very-high energy electrons, Klein-Nishina (KN) suppression of the IC process should be significant at multiple TeV energies causing a steepening in the VHE gamma-ray spectrum. For a target stellar photon field with mean temperature of 4500 K (Venter et al., 2009) and electron energies of 10 TeV the KN suppression factor is already about 0.025 (Coppi & Blandford, 1990), so the VHE gamma-ray spectrum should steepen well before this energy. In comparison, the observed spectrum may follow a straight power law, but this result is influenced by limited statistics. To account for the aforementioned arguments, as an alternative to an IC scenario, hadronic gamma-ray production is explored in this paper as the origin of the VHE source.

3. Ultrarelativistic outflow

3.1. Energetics

GRBs are generally believed to be caused by a pair of ultrarelativistic jets that are ejected from the central engine. Relativistic shock waves accelerate all particles from the incoming plasma to relativistic energies (Blandford & McKee, 1976), thus a substantial fraction of the initial energy of the relativistic blast wave is transferred into cosmic rays (see Atoyan et al., 2006, for the case of a GRB remnant). Therefore, for such a scenario, the energy in cosmic rays is a measure for the kinetic energy of the relativistic outflow.

From the luminosity of the VHE gamma-ray source the energy in hadronic cosmic rays can be estimated if the density of target material is assumed. At the location of Terzan 5 the density of target material should be on the order of $n \approx 0.1 \text{ cm}^{-3}$ (Clapson et al., 2011). Additionally, to constrain the total energy in hadronic cosmic rays for the entire relevant cosmic ray energy range above about 1 GeV a spectral index for the region below the range that can be probed with H.E.S.S. has to be assumed. If a cosmic ray spectral index of 2.0 is adopted below 5 TeV, the cosmic ray energy that produces gamma rays at the energy threshold of 440 GeV of the H.E.S.S. measurements, the total energy in hadronic cosmic rays would be $E_{\text{CR}} \approx 10^{51}(n/0.1\text{cm}^{-3})^{-1} \text{ ergs}$ (Abramowski et al., 2011).

Particles accelerated by a relativistic shock wave will not feature a single power-law spectrum but should have a break towards low energies where the break energy E_{br} is given by the bulk Lorenz factor Γ of the relativistic shock (Blandford & McKee, 1976; Katz, 1994) according to $E_{\text{br}} \sim m_p c^2 \Gamma^2 / 2$ with m_p the mass of the particle and c the speed of light. If a proton energy of 5 TeV, which produces gamma rays at the H.E.S.S. threshold, is adopted as the break energy, then the Lorenz factor at the time when most particles are accelerated would be $\lesssim 100$. Lower break energies would result in lower Lorenz factors. A spectral break towards lower particle energies would reduce the E_{CR} needed to explain the VHE source depending on the break energy by up to a factor of 2 (Atoyan et al., 2006). Since ultrarelativistic blast waves are expected to transfer a substantial part of their kinetic energy into cosmic rays the energetics of short GRBs could be roughly adequate for the observed VHE source if a ratio of prompt electromagnetic energy release to kinetic energy of 0.1 - 0.01 (Nakar, 2007) is assumed. Extending VHE observations to lower energies with the planned CTA array (CTA Consortium, 2010) could probe for a break in the gamma-ray spectrum of this source to test this scenario.

3.2. Age of the remnant and rate of short GRBs

After the acceleration, cosmic rays will diffuse away from the location of the GRB and will therefore form extended, center-filled gamma-ray sources (Atoyan et al., 2006). The age of a GRB remnant would in such a case be given by the diffusive propagation time of cosmic rays to the extension of the source. For the VHE source in Terzan 5 an age of the remnant of $10^3(D/10^{28} \text{ cm}^2 \text{s}^{-1})^{-1}$ years would be found (Abramowski et al., 2011), where D is the uncertain diffusion coefficient here compared to the value estimated for 5 TeV protons in the galactic disk of $10^{28} \text{ cm}^2 \text{s}^{-1}$ (Atoyan et al., 2006).

The age obtained for the potential GRB remnant at Terzan 5 can be compared to the rate of compact binary mergers in the Galaxy and to the rate of short GRBs in the local Universe. From field NSNS binaries a galactic merger rate to one event per $(0.5 - 7) \times 10^4$ years is found (Kalogera et al., 2004). For merger induced bursts which are formed in globular clusters a local rate of 20 - 90 events per $\text{Gpc}^{-3} \text{yr}^{-1}$ (Salvatera et al., 2008) or $\sim 4 \text{ Gpc}^{-3} \text{ yr}^{-1}$ (Guetta & Stella, 2009) has been estimated. With a density of Milky way-type galaxies in the local Universe of 0.01 galaxies per Mpc^{-3} (Cole et al., 2001) this results in a rate of short bursts per galaxy of about one event per $(0.1 - 0.5) \times 10^4 (f_b^{-1}/100)^{-1}$ years or $2.5 \times 10^4 (f_b^{-1}/100)^{-1}$ years. Here f_b is the beaming factor of short bursts, uncertain in the range of $1 \ll f_b^{-1} < 100$ (Nakar, 2007). It appears that the age of the potential GRB remnant would be roughly comparable to the rate of short bursts in the Milky Way if the diffusion coefficient D were below $10^{27} \text{ cm}^2 \text{s}^{-1}$ in the TeV range. Slowed down diffusion at

the location of Terzan 5 might be a reasonable assumption since X-ray observations demonstrated that the plasma there is highly turbulent (Yao & Wang, 2007; Crocker & Aharonian, 2011). To conclude, for possible diffusion coefficients the age estimated for the hypothetical remnant could be roughly comparable to the merger occurrence.

3.3. Potential multiwavelength signatures

Nonthermal diffuse X-ray emission from Terzan 5 extending beyond the half-mass radius has also been detected (Eger et al., 2010). The surface brightness of this X-ray emission roughly follows the surface brightness of the stellar component of the globular cluster, which could indicate an IC origin. In the framework of a GRB remnant scenario these diffuse X-rays might be emitted by primary electrons accelerated by the ultrarelativistic blast wave. Primary electrons would have cooled down below GeV energies since their production and therefore would not emit synchrotron radiation in the GHz band (Atoyan et al., 2006) but might still be energetic enough to scatter the intense stellar radiation field in the globular cluster to the X-ray range. For such a scenario an energy in electrons of $5 \times 10^{49} (u_{\text{rad}}/40 \text{ eV cm}^{-3})$ ergs would be required (Eger et al., 2010) with u_{rad} being the mean energy density in the stellar radiation field in the diffuse X-ray source. If this situation is indeed realized in Terzan 5, then a comparison to the energy in hadronic cosmic rays required to explain the VHE gamma-ray emission can be made. Here a quite large electron-to-proton ratio of ≈ 0.1 even after the electrons have cooled down below GeV energies would be found. Observations in the hard X-ray regime to extend the spectrum of the X-ray source might shed more light on its origin.

4. Subrelativistic ejecta

4.1. Pressure driven remnant

During the merger of compact binaries, a small fraction of the NS matter (10^{-4} - $0.1 M_{\odot}$) is dynamically ejected from the system with subrelativistic velocities (e.g. Ruffert et al., 1996; Lee & Klužniak, 1999; Rosswog, 2005). When assuming a velocity of the ejected material of $c/3$ (c speed of light), its kinetic energy would be 10^{49} - 10^{52} ergs. After the ejecta has displaced a comparable mass of ambient medium ($n \approx 0.1 \text{ cm}^{-3}$), most of the kinetic energy is converted into thermal energy (e.g. Dorfi, 1990) and the remnant expands farther in the pressure-driven phase (Domainko & Ruffert, 2005, 2008). It should, however, be noted that a large population of hadronic cosmic rays produced in the ultrarelativistic outflow may alter the evolution of the non-relativistic blast wave (e.g. Chugai et al., 2011). In the pressure-driven stage, the remnant is filled with hot thermal plasma that might be observable in thermal X-ray emission. Cooling of the thermal plasma appears to be unimportant in a remnant of the proposed age of a few times 10^4 years, since material with a density of 0.1 cm^{-3} and temperature of 10^7 (10^6) K would cool on a timescale of about 10^7 (3×10^5) years when X-ray line emission is considered (Sarazin, 1986). Not detecting an additional thermal component above the galactic diffuse X-ray emission and the extended nonthermal X-ray source with *Chandra* (Eger et al., 2010) would favor a content of hot thermal plasma in the potential remnant at the lower end of model predictions. A pressure driven remnant with energetics of 10^{49} ergs and an assumed age of 10^4 years would extend to a radius of about 10 pc corresponding to an angular radius of $6'$ at a distance of 5.9 kpc. It would be filled with thermal plasma of

about $T \approx 10^7$ K for a mean density of $n \approx 0.1 \text{ cm}^{-3}$, which should radiate $4 \times 10^{33} (T/10^7 \text{ K})^{-0.6} (n/0.1 \text{ cm}^{-3})^2 \text{ ergs s}^{-1}$ in X-rays, corresponding to $10^{-12} (T/10^7 \text{ K})^{-0.6} (n/0.1 \text{ cm}^{-3})^2 \text{ ergs cm}^{-2} \text{ s}^{-1}$ at Earth (X-ray emissivity from Sarazin, 1986). Such an extended and faint structure might have been missed by the present X-ray observations due to the limited field of view of *Chandra* and the limited sensitivity of *ROSAT*. Deep future X-ray observations with the existing *XMM* and the planned *IXO* mission (Barcons et al., 2011) would have the potential to probe the presence of diffuse thermal X-ray emission in the direction of Terzan 5.

4.2. Signatures from the ejecta

The ejecta of compact binary mergers may consist of r-process nuclei (e.g. Lattimer & Schramm, 1974; Ruffert et al., 1997; Freiburghaus et al., 1999). Many of these heavy nuclei are radioactive and emit hard X-ray and soft gamma-ray line emission during their decay (e.g. Qian et al., 1999; Domainko & Ruffert, 2005, 2008). The strength of these nuclear lines depend on the decay properties of the respective nuclei (Qian et al., 1999). Promising nuclei in terms of detectability of the gamma-ray line emission comprise a half-life time τ comparable to the age of the remnant. The estimated flux values in the table are given for a mass of the respective nuclei of $10^{-5} M_{\odot}$, an age of the remnant of 10^4 years and a distance to Terzan 5 of 5.9 kpc. Even for such an advantageous case, the line brightness would not exceed $10^{-8} \gamma \text{ cm}^{-2} \text{ s}^{-1}$, and detection would be challenging for the next generation of hard X-ray / soft gamma-ray instruments (e.g. Knöldlseder et al., 2009).

r-process nucleus	τ [10 ³ yr]	E_{γ} [keV]	F_{γ} [10 ⁻⁷ $\gamma \text{ cm}^{-2} \text{ s}^{-1}$]
²²⁶ Ra	2.31	609	0.01
²²⁹ Th	10.6	40.0	0.04
²⁴³ Am	7.37	74.7	0.09

Table 1. Properties of gamma ray lines in a potential merger remnant in Terzan 5.

5. Ionizing radiation

5.1. Astrophysical signatures

GRBs and their afterglows exhibit strong ionizing radiation which should alter the equilibrium state of the surrounding interstellar medium (ISM). Signatures of such an ionizing event should be observable for several 10^4 years in form of high-ionization lines (Perner et al., 2000). Terzan 5 is, however, located in a direction close to the galactic center ($l = 3.8^\circ$, $b = 1.7^\circ$), hence affected by severe galactic absorption ($A_V = 7.72$, Ortolani et al., 1996), which may challenge the detection of such emission lines.

5.2. Terrestrial signatures

The ionizing radiation of a GRB that is beamed towards the Earth may also alter the state of the Earth's atmosphere (Fishman & Inan, 1988), and Galactic GRBs should produce a substantial amount of NO_y there (Melott et al., 2004). This can be deposited as nitrates in the polar ice caps (e.g. Thomas et al.,

2005; Ejzak et al., 2007), which represent an archive of the atmospheric conditions for 7.4×10^5 years (Augustin et al., 2004). Similar signatures of nearby, Galactic supernovae have already been claimed (e.g. Motizuki et al., 2009). Here the possibility of such a signature originating from a GRB in Terzan 5 is appraised. Considerations in this section are less well founded than discussions in previous sections. An evaluation of different terrestrial NO_y signatures can be found in Harfoot et al. (2007).

To assess NO_y deposition from a GRB in Terzan 5, the solar proton event that occurred during August 1972 and which is imprinted in polar ice (Zeller et al., 1986; Shea et al., 2006) is used as a reference. This specific event produced 3.6×10^{33} molecules NO_y in the middle atmosphere (Jackman et al., 2005). It is found that a fluence of about 10^5 ergs/cm² of ionizing radiation is needed for a comparable NO_y enhancement if it is adopted that 35 eV are necessary per ionization (Gehrels et al., 2003) and that each ionization releases 1.25 N atoms, which rapidly produce NO_y (Jackman et al., 2005). A GRB at an adopted distance of Terzan 5 of 5.9 kpc would have to release an isotropic fluence comparable to typical short bursts of several 10^{50} ergs in ionizing radiation to produce enough nitrates to leave significant signatures in polar ice cores. Thus such an event could indeed have left traces in polar ice but only on the level of more frequent solar proton events. However, since GRBs are beamed with a beaming factor f_b in the range of $1 \ll f_b^{-1} < 100$, the probability that the prompt GRB emission has hit the Earth would only be a few percent.

Along with generating nitrates by X-rays and soft gamma-rays, photons from the GRB with energies > 10 MeV would produce long-living radionuclides like ¹⁰Be (half-life time 1.5 Myr) in the Earth atmosphere (Thorsett, 1995), which will after rain-out also be deposited in polar ice caps. Such high-energy emission has indeed been observed from some short GRBs (e.g. Abdo et al., 2010b). In this framework it is interesting to note that two possibly cosmogenic ¹⁰Be anomalies have been found in antarctic ice cores (Raisbeck et al., 1987). However, the detectability of such signatures from a potential short GRB in Terzan 5 will again be limited by the beaming of the high-energy radiation.

6. Summary

The VHE gamma-ray source discovered in the vicinity of Terzan 5 appears to be compatible with being the remnant of a short GRB induced by the merger of two compact stars. An incontrovertible proof is missing, but also no contradictory facts have been found for this scenario. Further observations in the X-ray and gamma-ray regime may help to strengthen the evidence for a GRB remnant in Terzan 5. In particular X-ray observations could probe the presence of thermal plasma heated by subrelativistic merger ejecta. The identification of a GRB remnant in our own Galaxy would provide important constraints on the GRB rate in the very local universe and may also help to explore the effect of such a highly energetic event on the Earth.

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References

- Abdo, A. A. et al. (The Fermi LAT collaboration) 2010a, A&A, 524, 75
 Abdo, A. A. et al. (The Fermi LAT collaboration) 2010b, ApJ, 716, 1178
 Abramowski, A. et al. (The H.E.S.S. collaboration) 2011, A&A accepted, arXiv:1106.4069
 Aharonian, F. (The H.E.S.S. collaboration) 2005, A&A, 439, 1013
 Aharonian, F. (The H.E.S.S. collaboration) 2008, A&A, 477, 353
 Anderson, S. B., Gorham, P. W., Kulkarni, S. R., Prince, T. A., & Wolszczan, A. 1990, Nature 346, 42
 Atoyan, A., Buckley, J. & Krawczynski, H. 2006, ApJ, 642, L153
 Augustin, L., Barbante, C., Barnes, P. R. et al. 2004, Nature, 429, 623
 Barcons, X. et al. 2011, arXiv:1102.2845
 Bednarek, W. & Sitarek, J. 2007, MNRAS, 377, 920
 Blandford, R. D., & McKee, C. F. 1976, Phys. Fluids, 19, 1130
 Chugai, N. N., Churazov, E. M. & Sunyaev, R. A. 2011 arXiv:1101.5291
 Clapson, A.-C., et al. 2011, A&A accepted, arXiv:1105.5579
 Cole, S., Norberg, P., Baugh, C. M. et al. 2001, MNRAS, 326, 255
 Coppi, P. S. & Blandford, R. D. 1990, MNRAS, 245, 453
 Crocker, R., M. & Aharonian, F. A. 2011, PRL, 106, 101102
 CTA Consortium 2010, arXiv:1008.3703
 Domainko, W. & Ruffert, M. 2005, A&A, 444, L33
 Domainko, W. & Ruffert, M. 2008, AdSpR, 41, 518
 Dorfi, E.A. 1990, A&A, 234, 419
 Eger, P., Domainko, W. & Clapson, A.-C. 2010, A&A, 513, 66
 Ejzak, L. M., et al. 2007, ApJ, 654, 373
 Ferraro, F. R., Beccari, G., Dalessandro, E., et al. 2009, Nature, 462, 1028
 Fishman, G. J. & Inan, U. S. 1988, Natur, 331, 418
 Freiburghaus, C., Rosswog, S., & Thielemann, F.-K. 1999, ApJ, 525, 121
 Gehrels, N., Laird, C. M., Jackman, C. H., Cannizzo, J. K., Mattson, B. J. & Chen, W. 2003, ApJ, 585, 1169
 Gehrels, N., Sarazin, C. L., O'Brien, P. T., et al. 2005, Natur, 437, 851
 Gehrels, N., Ramirez-Ruiz, E. & Fox, D. B. 2009, ARA&A, 47, 567
 Grindlay, J., Portegies Zwart, S., & McMillan, S. 2006, NatPh, 2, 116
 Guetta, D. & Stella, L. 2009, A&A, 498, 329
 Harfoot, M. B. J., Beerling, D. J., Lomax, B. H. & Pyle, J. A. 2007, JGR, 112, D07308
 Ioka, K., Kobayashi, S. & Mészáros, P. 2004, ApJ, 613 L17
 Ioka, K. & Mészáros, P. 2010, ApJ, 709, 1337
 Jackman, C. H., Deland, M. T. Labow, G. J. et al. 2005, AdSpR, 35, 445
 Kalogera, V., Kim, C., Lorimer, D. R., et al. 2004, ApJ, 601, L179, erratum 2004, ApJ, 614, L137
 Katz, J. I. 1994, ApJ, 432, L107
 Knöldlseder, J., von Ballmoos, P., Frontera, F., et al. 2009, ExA, 23, 121
 Kong, A. K. H., Hui, C. Y. & Cheng, K. S. ApJ, 712, L36
 Lanzoni, B., Ferraro, F. R., Dalessandro, E., et al. 2010, ApJ, 717, 653
 Lattimer, J. M. & Schramm, D. N. 1974, ApJ, 192, L145
 Lee, W. H. & Kluźniak, W. 1999, ApJ, 526, 178
 Lee, W. H., Ramirez-Ruiz, E., & Granot, J. 2005, ApJ, 630, L165
 Melott, A. L., Lieberman, B. S., Laird, C. M., et al. 2004, IJAsB, 3, 55
 Motizuki, Y., Takahashi, K., Makishima, K., et al. 2009, submitted to Nature, arXiv:0902.3446v1
 Nakar, E. 2007, PhR, 442, 166
 Ortolani, S., Barbuy, B. & Bica, E. 1996, A&A, 308, 733
 Perner, R., Raymons, J. & Loeb, A., 2000, ApJ, 533, 658
 Pooley, D. & Hut, P. 2006, ApJ, 646, L143
 Qian, Y.-Z., Vogel, P. & Wasserburg, G. J. 1999, ApJ, 524, 213
 Raisbeck, G. M.; Yiou, F.; Bourles, D.; Lorian, C.; Jouzel, J.; Barkov, N. I. 1987, Natur, 326, 273
 Ransom, S. M. 2008, AIPC, 983, 415
 Rosswog, S. 2005, ApJ, 634, 1202
 Ruffert, M., Janka, H.-T. & Schaefer, G. 1996, A&A, 311, 532
 Ruffert, M., Janka, H.-T., Takahashi, K. & Schaefer, G. 1997, A&A, 319, 122
 Salvaterra, R., Cerutti, A., Chincarini, G., Colpi, M., Guidorzi, C. & Romano, P. 2008, MNRAS, 388, L6
 Sarazin, C. L. 1986, RvMP, 58, 1
 Shea, M. A., Smart, D. F., McCracken, K. G., Dreschhoff, G. A. M. & Spence, H. E. 2006, AdSpR, 38, 232
 Thomas, B. C., Melott, A. L., Jackman, C. H., et al. 2005, ApJ, 634, 509
 Thorsett, S. E. 1995, ApJ, 444, L35
 Venter, C., de Jager, O. C. & Clapson, A.-C. 2009, ApJ, 696, L52
 Yao, Y. & Wang, Q. D. 2007, ApJ, 666, 242
 Zeller, E. J., Dreschhoff, G. A. M. & Laird, C. M. 1986, GeoRL, 13, 1264